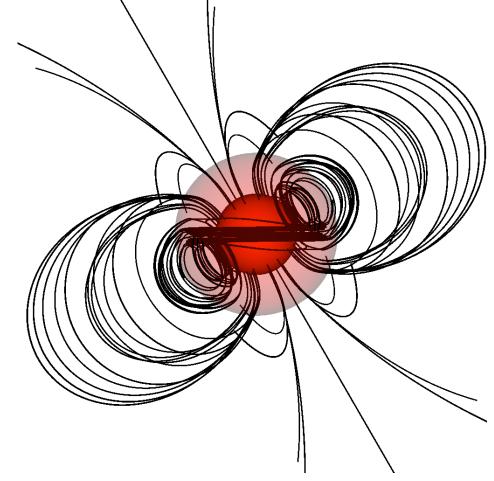




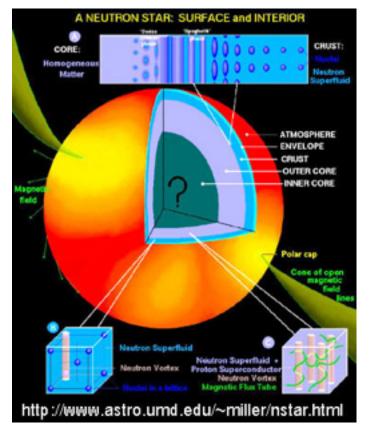
Determining observational properties of neutron stars by modelling their interiors

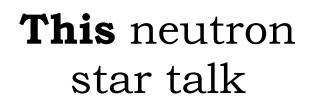
Paul Lasky

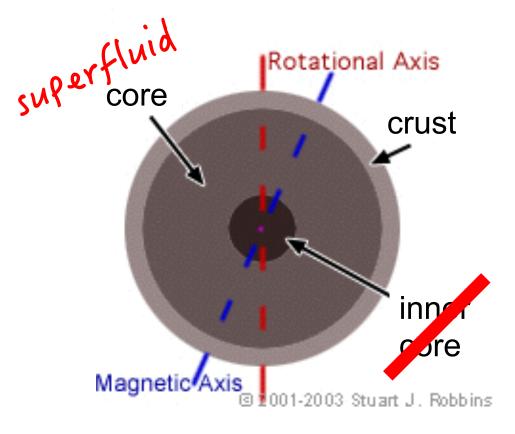


Neutron Star Anatomy

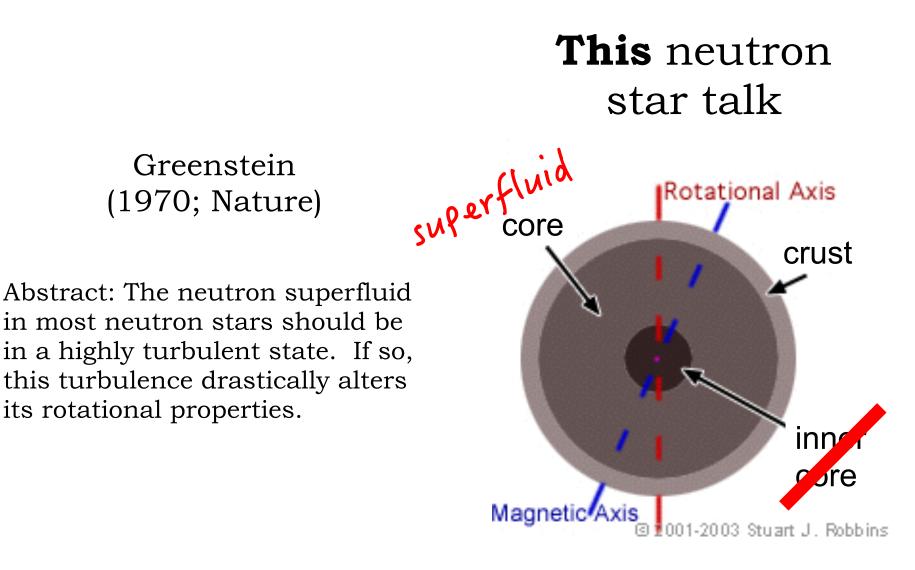
Fiducial neutron star talk

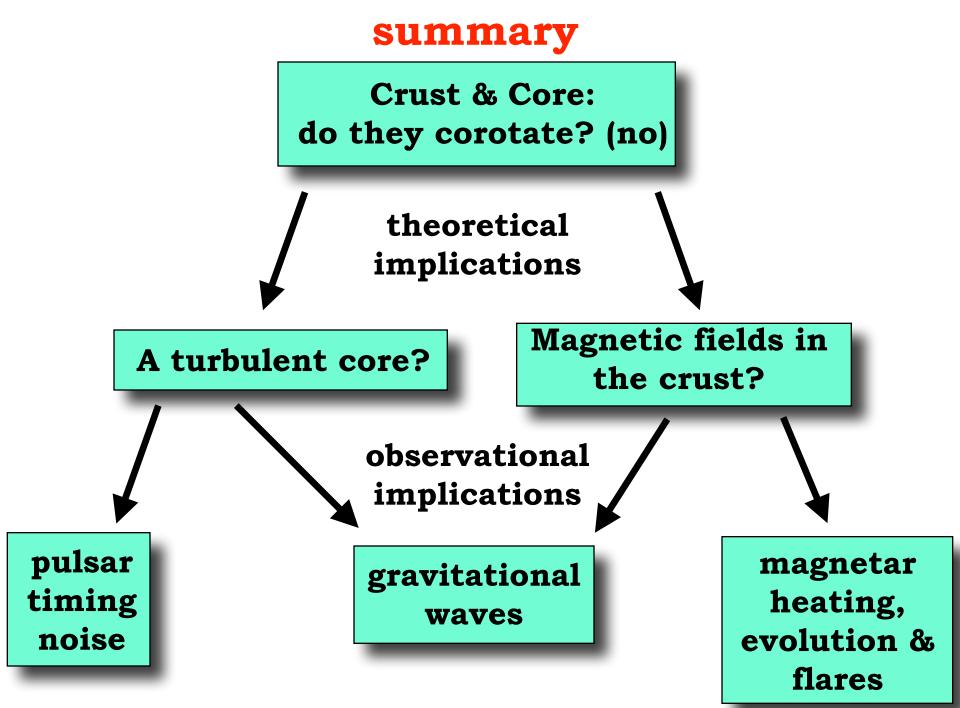






Neutron Star Anatomy





The core and the crust

Conventional wisdom:

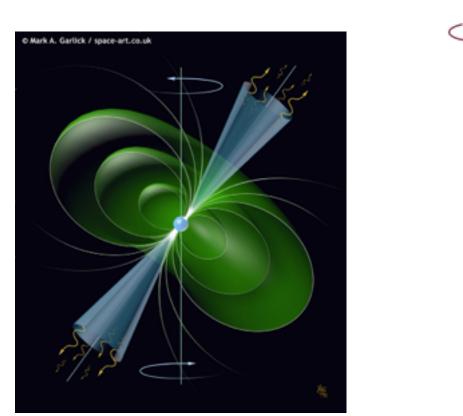
- Neutron star's crust & core *corotate*
- 2 mechanisms:
 - viscous coupling (Ekman pumping)
 - magnetic coupling (commonly considered dominant)

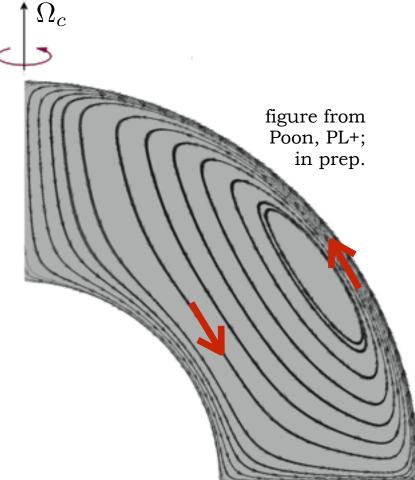
The conventional wisdom is wrong!

Neither mechanism can effectively enforce crustcore corotation (Melatos 2012; Glampedakis & PL 2015)

• Theoretical & Observational implications

Ekman pumping





magnetic field spins down crust

Ekman pumping spins down fluid in core

Ekman pumping

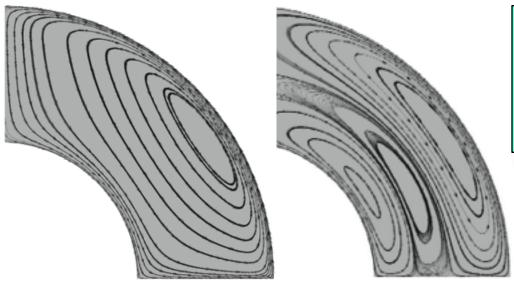


Vortex flow at the Canberra airport

"this is something that only physicists can get excited about..." Alexander Heger

stratified Ekman pumping

- Ekman flow hindered by stratification (Abney & Epstein 1996)
- Only effective in thin layer near crust-core boundary
- Rest of core couples on much longer timescale (~ 10³ yr; Melatos 2012)



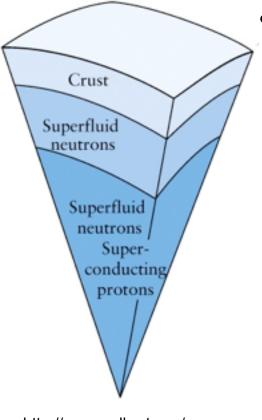
Melatos 2012: neutron stars have super-rotating cores!

caveat: the magnetic field!

figures from Poon, PL+; in prep.

Model

• Two-fluid core (charged proton-electron fluid + neutron superfluid) magnetically coupled to the crust.



http://www.ualberta.ca/

• in crust's instantaneous rest frame, the secular dynamics of charged component is

Glampedakis & PL (2015)

$$2\mathbf{\Omega} \times \mathbf{v}_{\mathbf{p}} + \dot{\mathbf{\Omega}} \times \mathbf{r} + \nabla \Psi_{p} = \frac{1}{\rho_{p}} \left(\mathbf{F}_{\text{mag}} - \mathbf{F}_{\text{cpl}} \right)$$
$$2\mathbf{\Omega} \times \mathbf{v}_{n} + \dot{\mathbf{\Omega}} \times \mathbf{r} + \nabla \Psi_{n} = \frac{1}{\rho_{n}} \mathbf{F}_{\text{cpl}}$$

 Ψ : chemical + gravitational potentials \mathbf{F}_{mag} : magnetic force \mathbf{F}_{cpl} : coupling force with neutrons

Glampedakis & PL (2015)

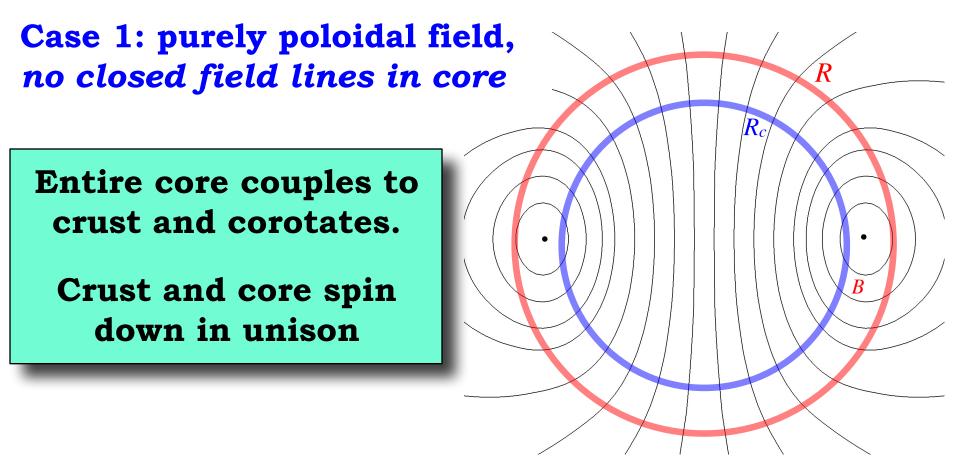
The punch line

• Degree of coupling between the crust and the core depends sensitively on the magnetic field **geometry!**

Glampedakis & PL (2015)

The punch line

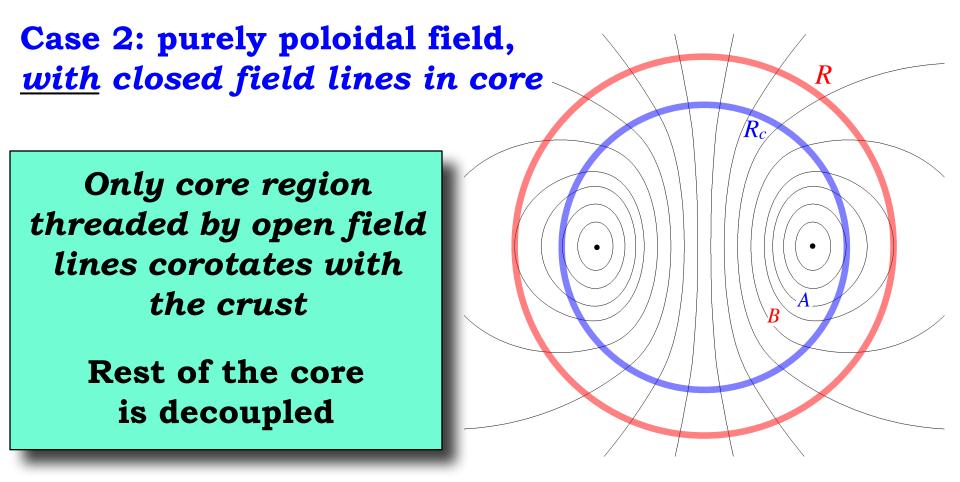
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Glampedakis & PL (2015)

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Glampedakis & PL (2015)

The punch line

• Degree of coupling between the crust and the core depends sensitively on the magnetic field **geometry!**

Case 3: mixed toroidal-poloidal field, with closed field lines in core Only core region threaded by open field lines corotates with the crust

Rest of the core is decoupled

the super-rotating core region

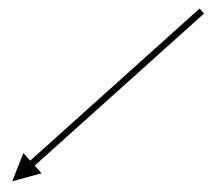
Glampedakis & PL (2015)

R

• Following birth, neutron stars could have a **super-rotating**, **torus-shaped region in the core!**

- Almost certainly unstable:
 - velocity jump along field line A induces local Lorentz force that will try to displace the super-rotating region
 - also should be unstable to Kelvin-Helmholtz instability

Two Possible Outcomes

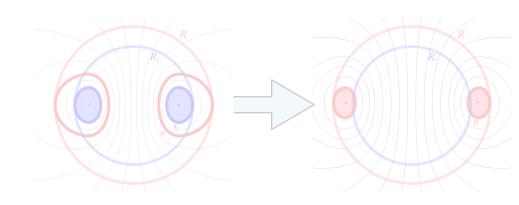


i) core remains in constant turbulent state



Peralta, Melatos, et al.

ii) magnetic field evicted to crust



Turbulent Consequences

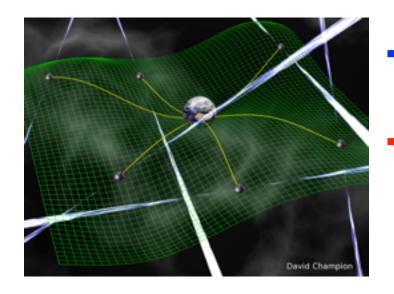
Pulsar timing noise

- ➡ Is pulsar timing noise from turbulence?
- Quantifying the effect on gravitational wave detection with Pulsar Timing Arrays

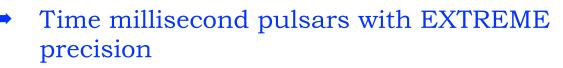
➡ Gravitational waves - LIGO

- Single Neutron Star
- Stochastic Background

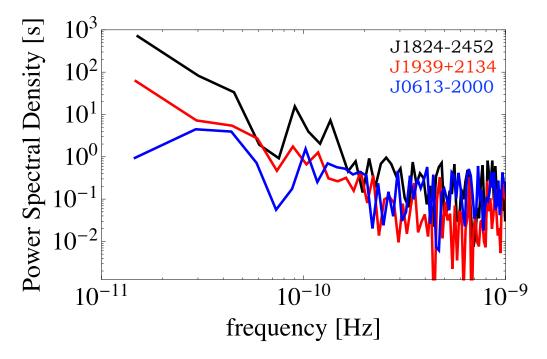
Pulsar Timing Arrays



- ➡ Timing noise:
 - Stochastic wandering of pulse arrival times
 - stochastic torque fluctuations exerted on crust by turbulent fluid in core



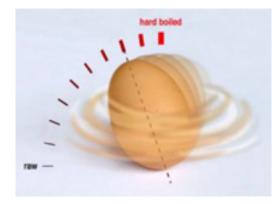
- Look for correlated timing residuals as GW signature
 - Timing Residual: difference between measured and modelled phase of pulse.



Timing Noise Due to Turbulence

Greenstein (1970; Nature)

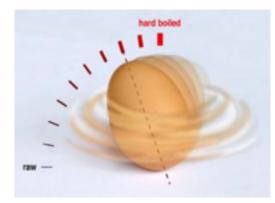
'My final point is a speculative one. When an uncooked egg rotates it does so irregularly. The yolk inside moves about erratically, and in order to conserve angular momentum the rotation rate of the shell must also fluctuate. The rotating turbulent neutron superfluid must exhibit something like the same phenomenon.



Timing Noise Due to Turbulence

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'My final point is a speculative one. When an uncooked egg rotates it does so irregularly. The yolk inside moves about erratically, and in order to conserve angular momentum the rotation rate of the shell must also fluctuate. The rotating turbulent neutron superfluid must exhibit something like the same phenomenon.

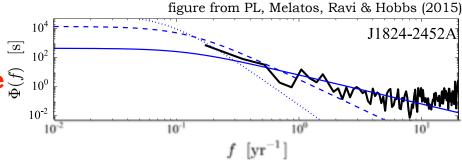


Melatos & Link (2014)

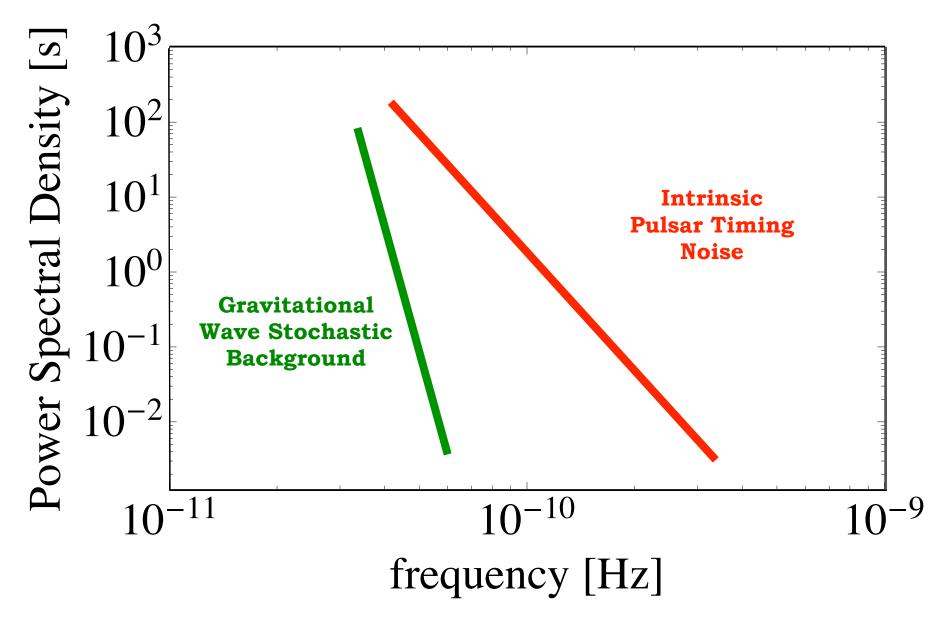
calculated angular momentum \square fluctuations on NS crust from core $\sum_{n=1}^{\infty}$

Small number of (relatively unconstrained) parameters 'fit' timing noise spectra.

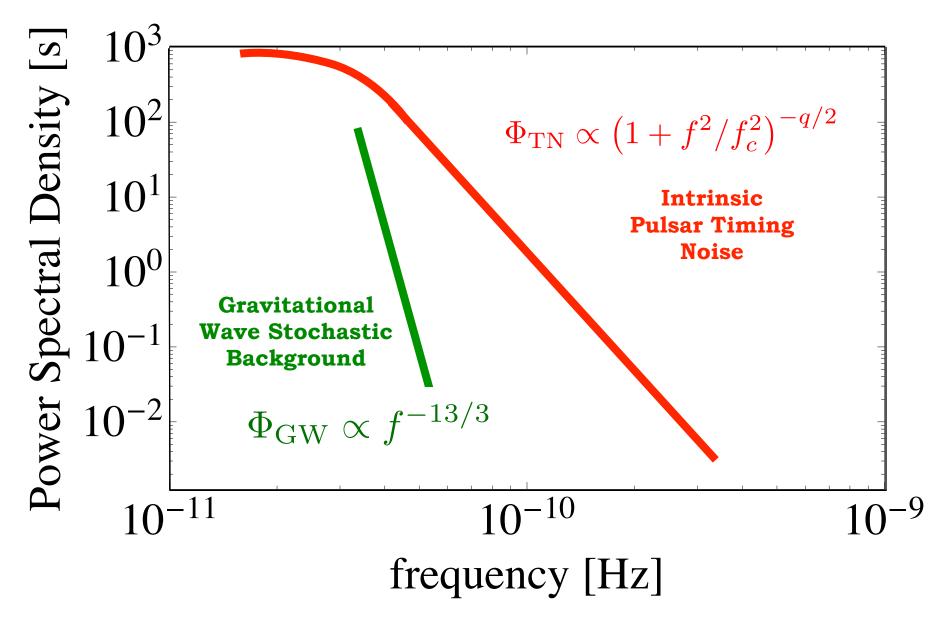
$f [yr^{-1}]$ Prediction: low-frequency plateau in timing noise spectrum



PL, Melatos, Ravi & Hobbs (2015)

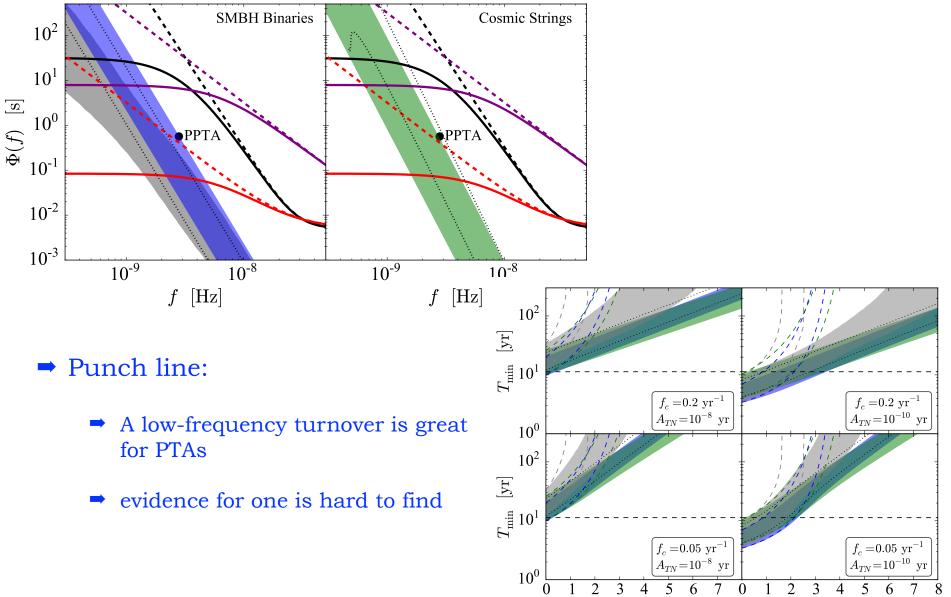


PL, Melatos, Ravi & Hobbs (2015)



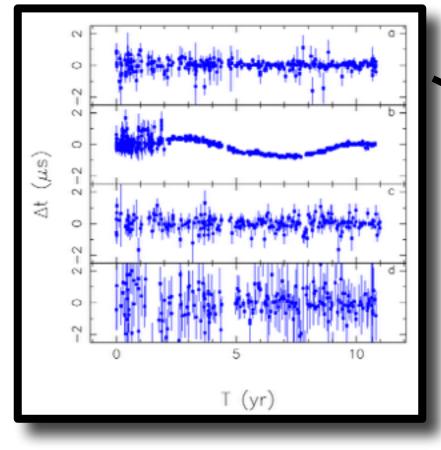
PL, Melatos, Ravi & Hobbs (2015)

q

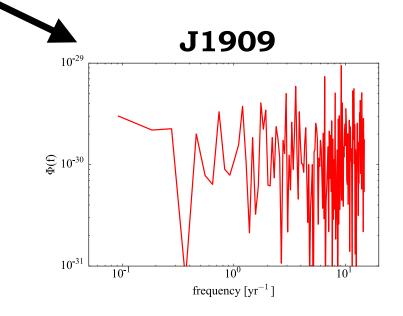


q

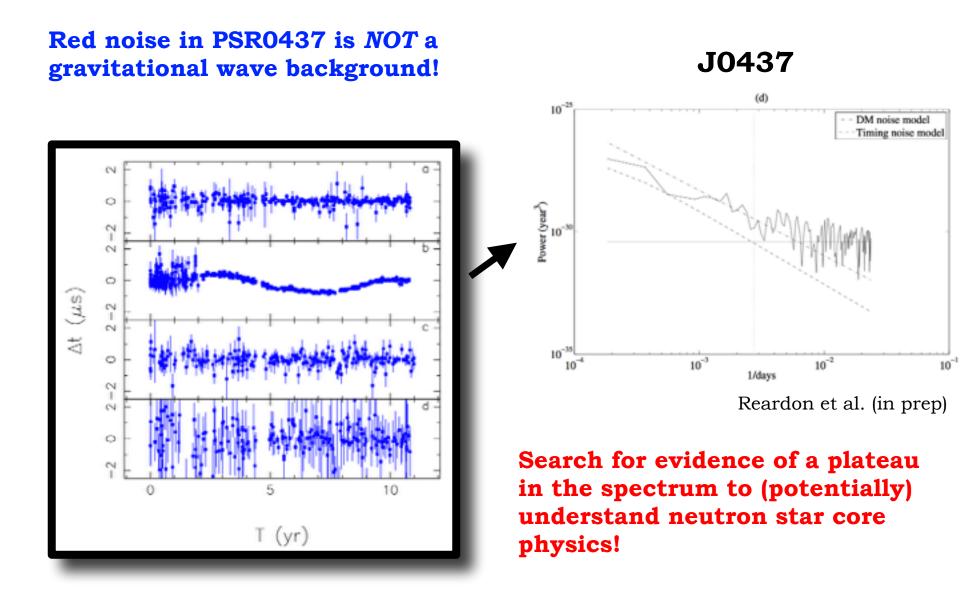




Bayesian analysis: No evidence for red noise in the two pulsars that are biggest contributors to new GW limit



What about the level of red noise?



Turbulent Consequences

➡ Pulsar timing noise

- ➡ Is pulsar timing noise from turbulence?
- Quantifying the effect it has on gravitational wave detection with PTAs

Gravitational waves - LIGO

- ➡ Single Neutron Star
- Stochastic Background

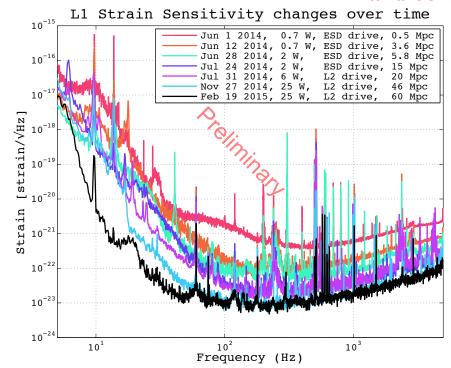
Turbulent flows emit gravitational waves

Melatos & Peralta (2010)

Turbulence: On average - axisymmetric Instantaneously non-axisymmetric

$$h_{\rm rms} = 5 \times 10^{-28} \left(\frac{M_{\star}}{1.4 \, M_{\odot}}\right) \left(\frac{R_{\star}}{10 \, \rm km}\right)^3 \left(\frac{d}{1 \, \rm kpc}\right)^{-1} \left(\frac{\Delta \Omega}{10 \, \rm rad \, s^{-1}}\right)^3 \underbrace{ \text{Shear}}_{\substack{\text{velocity between crust}\\ \text{and core}}}$$

Only of potential interest to LIGO for nearby (d ~10 pc) fast rotators (P ~ 1 ms $=> \Delta \Omega ~ 30$)



One neutron star does not emit a detectable gravitational wave signal What about ALL the neutron stars in the Universe?

PL, Bennett & Melatos (2013)

Consider 2 Populations

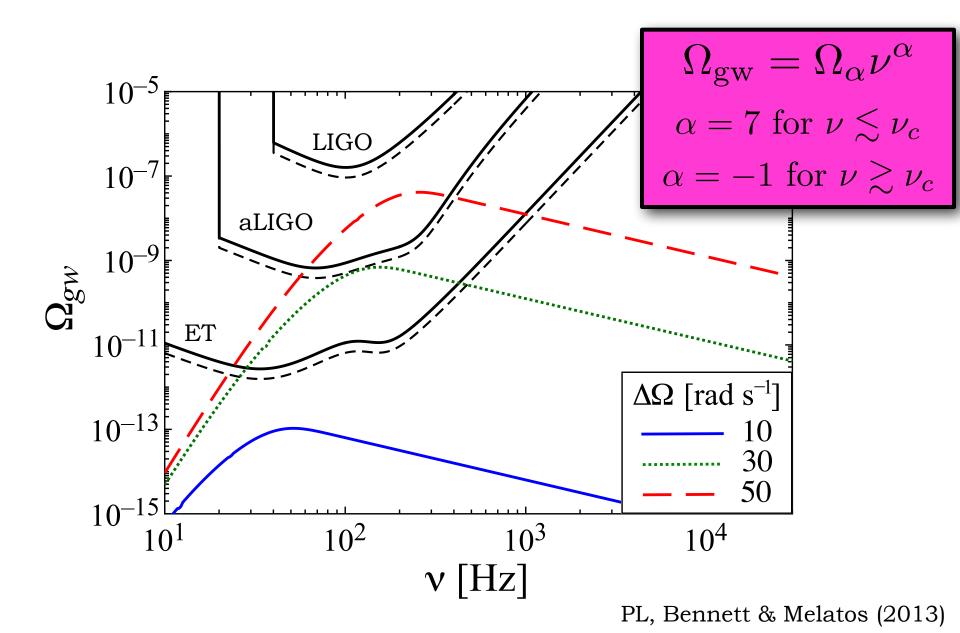
➡ Naive:

➡ all NSs in Universe have same ΔΩ. i.e. ΔΩ is independent of Ω

Radio pulsars:

Broad distribution of $\Delta \Omega$, where $\Delta \Omega$ is proportional to spindown rate

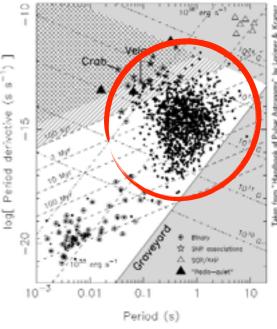
Unique $\Delta \Omega$



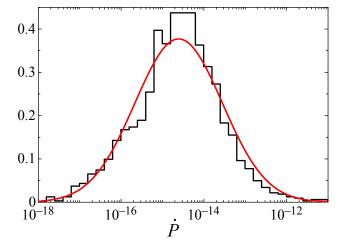
Pulsar Population

take galactic distribution of known pulsars and assume the

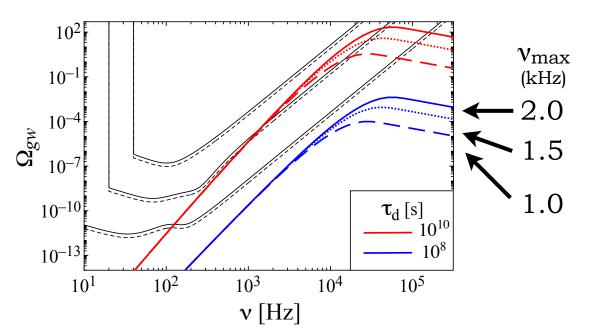




atnf.csiro.au



 $\Delta \Omega = \tau_d \Omega$



PL, Bennett & Melatos (2013)

Detection unlikely

"Things sometimes happen to me that are very, very unlikely. Does it mean that I don't exist with high Bayesian probability?" (Levin 2015, Facebook)

Non-detections give interesting constraints on shear damping times, etc.

$$\Omega_{gw} = \Omega_{\alpha} \nu^{\alpha}$$

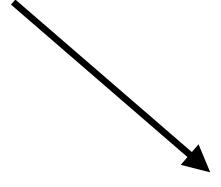
$$\alpha = 7 \text{ for } \nu \lesssim \nu_{c}$$

$$\alpha = -1 \text{ for } \nu \gtrsim \nu_{c}$$

It is worth searching
for this.

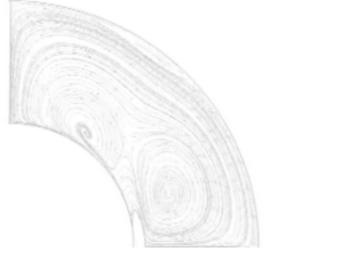
PL, Bennett & Melatos (2013)

Two Possible Outcomes

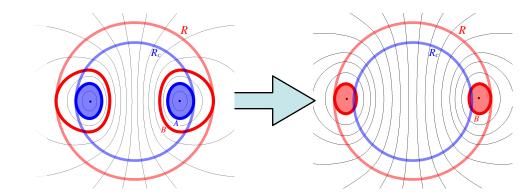


i) core remains in constant turbulent state

ii) magnetic field evicted to crust



Peralta, Melatos, et al.

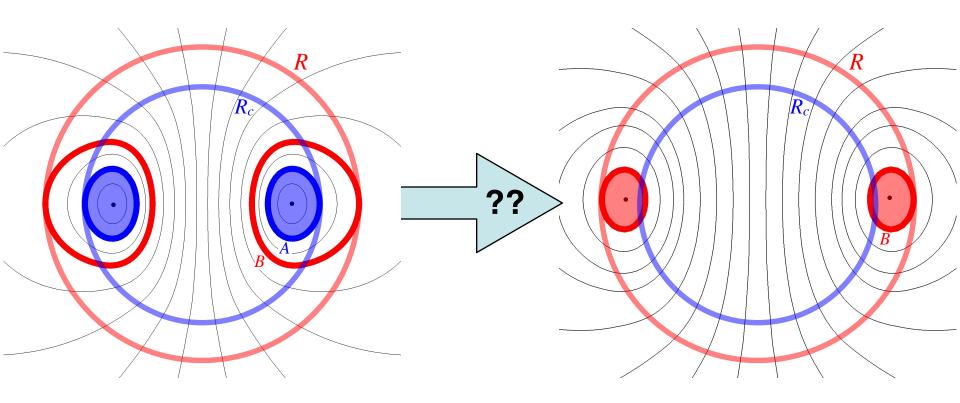


the crust as a magnetic field depository

Glampedakis & PL (2015)

• A Conjecture:

the system will evict the closed field lines + toroidal region into the crust

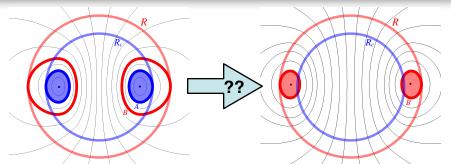


the crust as a magnetic field depository

Glampedakis & PL (2015)

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• young magnetars:

 $B\gtrsim 10^{15}\,{
m G}$: star spins down before crust forms (~1 day) $B\lesssim 10^{15}\,{
m G}$: our model applies

hydrodynamic instability timescale

$$t \sim t_{\rm sd} \approx 4.7 \left(\frac{B_p}{10^{15}\,{\rm G}}\right)^{-2} \left(\frac{P}{10\,{\rm ms}}\right)^2\,{\rm d}$$

 $t \lesssim t_{sd}$: gravitational wave emission $t \gg t_{sd}$: eviction of the magnetic field to the core

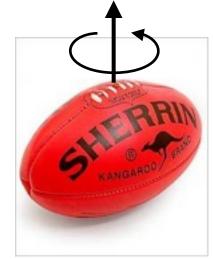
$t \lesssim t_{ m sd}$: gravitational wave emission

• Strong toroidal field wound up in core

$$h_0 \propto \frac{I\epsilon\nu^2}{D}$$

• ε due to magnetic deformations

$$\epsilon \sim 10^{-6} \left(\frac{B_{\rm t}}{10^{15} \, {\rm G}}
ight)$$
 Haskell et al. (2008, erratum 2009),
Mastrano et al. (2011)



$\lesssim t_{\rm sd}$: gravitational wave emission

Strong toroidal field wound up in core

$$h_0 \propto \frac{I\epsilon\nu^2}{D}$$

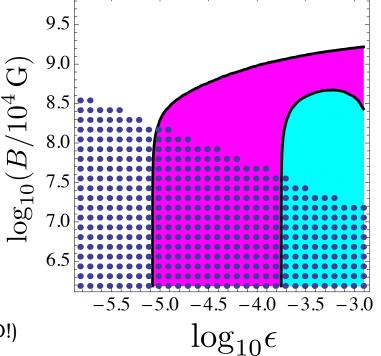
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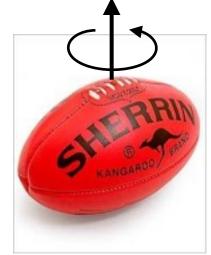
$$\epsilon \sim 10^{-6} \left(\frac{B_{\rm t}}{10^{15} \, {\rm G}} \right) \overset{\rm e.g., \ {\rm Cutler} \ (2002)}{\rm Haskell \ et \ al.} \ (2008, \ \frac{\rm erratum}{\rm 2009}), \\ {\rm Mastrano \ et \ al.} \ (2011)$$

unfortunately, need long integration times:

SN1987A: ~ 1 yr integration Chung, Melatos, et al. (2011)

(important: other emission mechanisms possible that make this search worthwhile — Lilli Sun's PhD!)







• Strong toroidal field wound ι

$$h_0 \propto rac{I\epsilon
u^2}{D}$$

ε due to magnetic deformatic

$$\epsilon \sim 10^{-6} \left(\frac{B_{\rm t}}{10^{15} \, {\rm G}} \right)$$
 Haskell

$$\epsilon \sim 10^{-6} \left(\frac{B_t}{10^{15} \,\mathrm{G}} \right)^{\mathrm{Haskell}}$$
 $\epsilon^{CFL} \sim 2.5 \times 10^{-4} \left(\frac{\langle B_t \rangle}{10^{15} \,\mathrm{G}} \right)$
unfortunately, need lon
integration times: This is potentially a nuclear
physics experiment!

otentially a nuclear physics experiment!

Unless....

 $\epsilon^{2SC} \sim 8.0 \times 10^{-5} \left(\frac{\langle B_t \rangle}{10^{15} \,\mathrm{G}} \right)$

Owen (2004), Glampedakis et al. (2012)

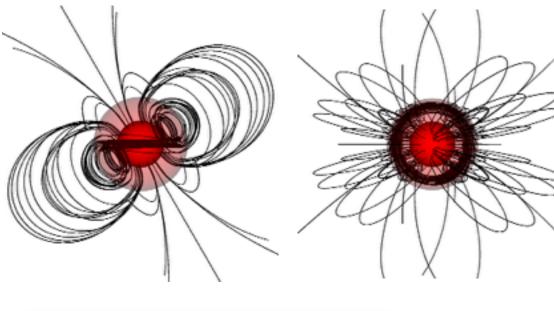
SN1987A: ~ 1 yr integrati Chung, Melatos, et al. (2011)

(important: other emission mechanisms possible that make this search worthwhile — Lilli Sun's PhD!

$$\begin{array}{c} 7.5 \\ 0 \\ 0 \\ 0 \\ -5.5 \\ -5.5 \\ -5.0 \\ -4.5 \\ -4.0 \\ -3.5 \\ -3.0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 10 \\ \epsilon \end{array} \right)$$

$\lesssim t_{ m sd}$: gravitational wave emission

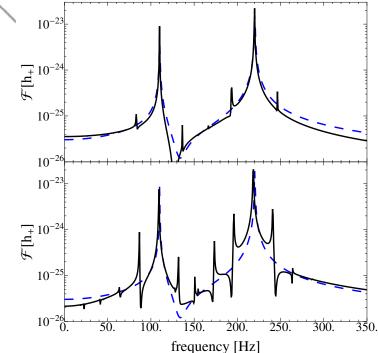
• a positive detection also allows us to probe the stellar geometry (e.g., Mastrano, PL & Melatos 2013 for multipolar fields)



e.g., twisted-torus

PL & Melatos (2013)

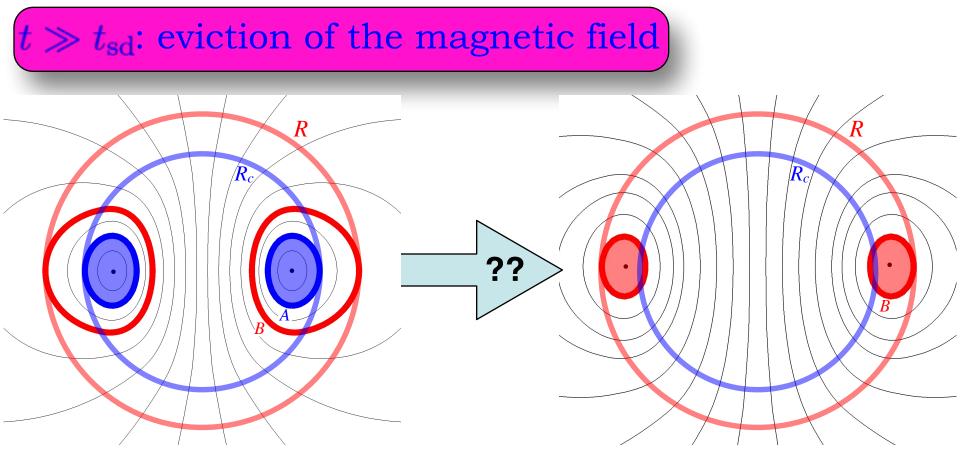
- triaxial deformation
- 'naturally motivated'
- Enriches GW signal



e.g., NS born with:

- Virgo cluster (~20 Mpc)
- $B_p = 10^{14} G$
- $\langle B_t \rangle = 10^{16} \text{ G}$

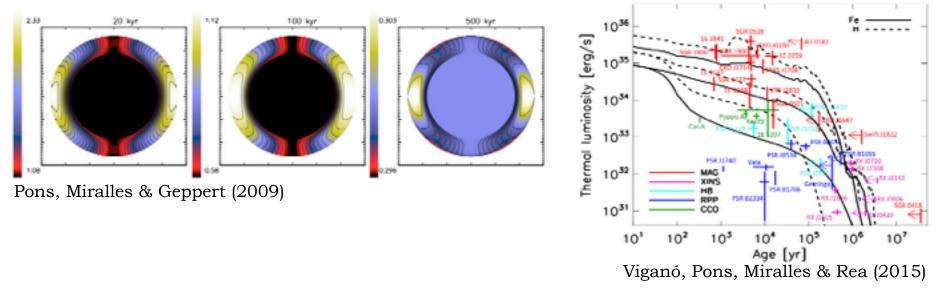
Triaxiality must last ~ one month for SNR~3 in aLIGO



✓existence of strong toroidal field in crust is key for magnetar heating, fast magnetic evolution and flares!

$t \gg t_{ m sd}$: eviction of the magnetic field

- magnetar heating
 - e.g., series of papers by Pons & collaborators, Ho et al. 2012
 - magneto-thermal evolution of strong crustal fields



- magnetar flares
 - giant flares is crust-fracturing by strong B-field involved? (e.g., Thompson & Duncan series)

the future?

magnetic field does not couple the core and crust of a neutron star.

Conjecture: stability is reached when closed field lines + toroidal field are evicted into crust.

- what's next?
 - more general B-field geometry
 - easy to generalise to higherorder multipoles
 - non-axisymmetric more difficult!
 - superconducting MHD

• how does the system *actually* evolve?

